

Modification and implementation of the CABARET scheme in the Coupled Climate Model INM RAS

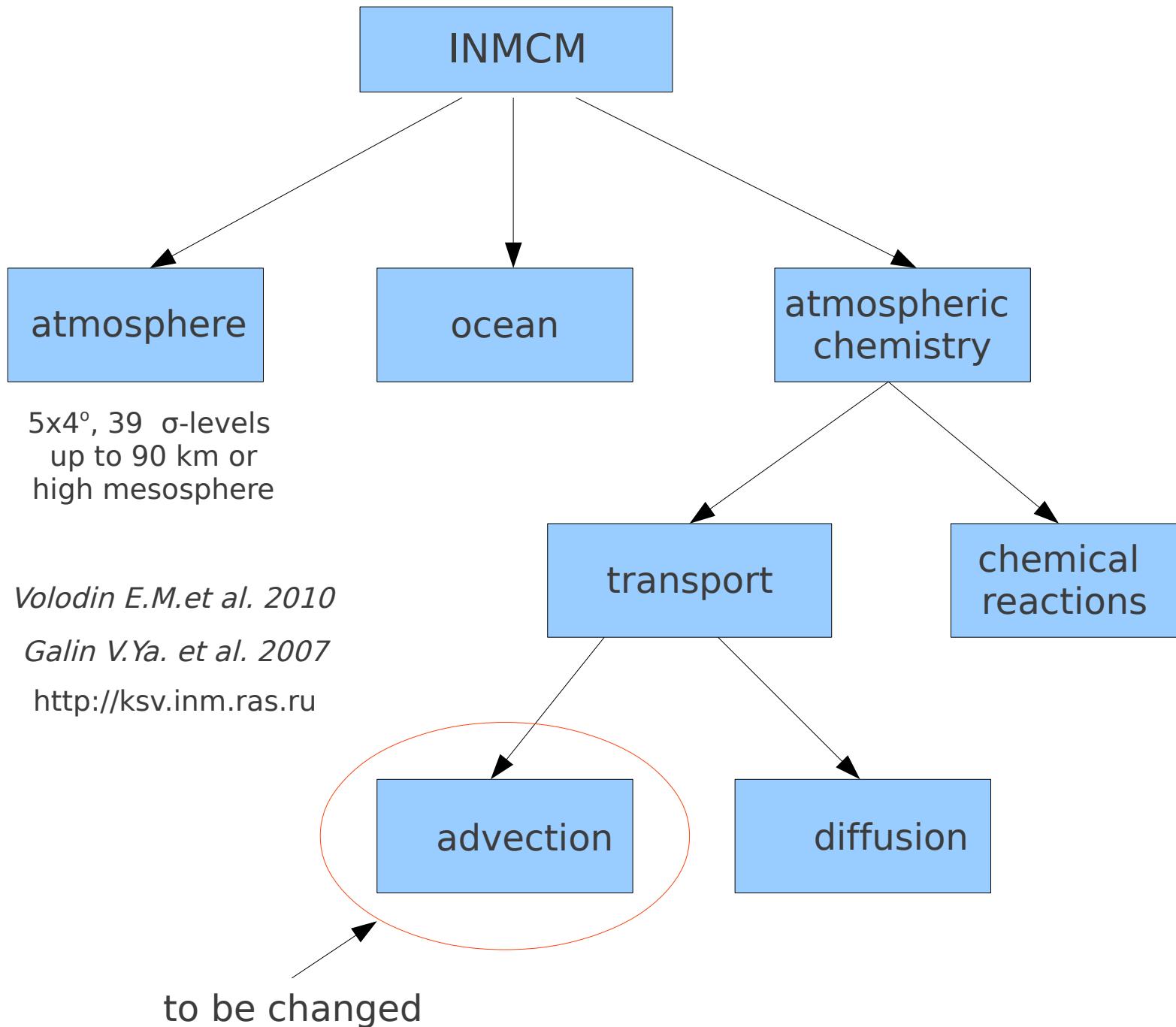
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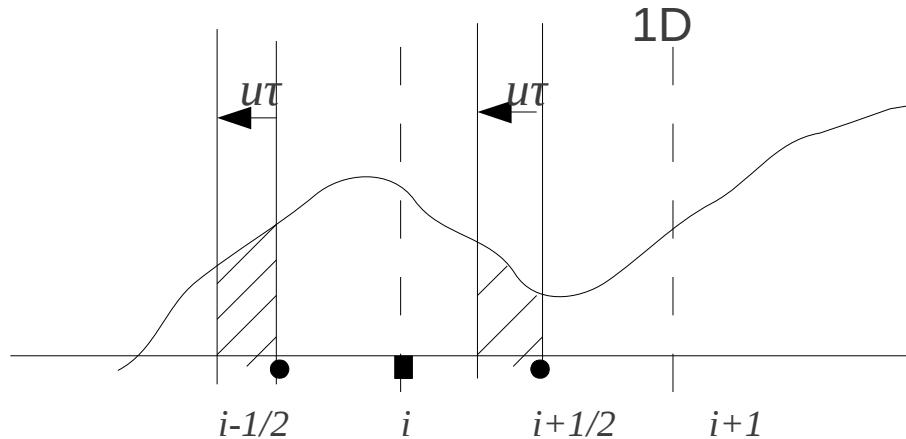
Coupled Climate Model of INM RAS



CABARET advection scheme in 1D case

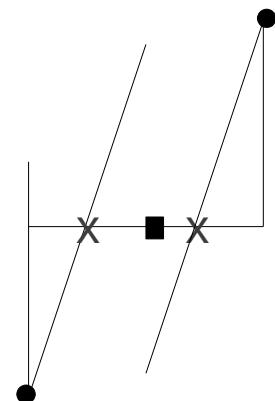
CABARET - Compact Accurately Boundary-Adjusting high-REsolution Technique
(upwind leapfrog)

Goloviznin et al., 2003



**Flux form equation
for conservative variables**

$$\frac{\rho_i^{n+1} - \rho_i^n}{\tau} + \frac{(u\rho)_{i+\frac{1}{2}}^{n+\frac{1}{2}} - (u\rho)_{i-\frac{1}{2}}^{n+\frac{1}{2}}}{\Delta x_i} = 0$$

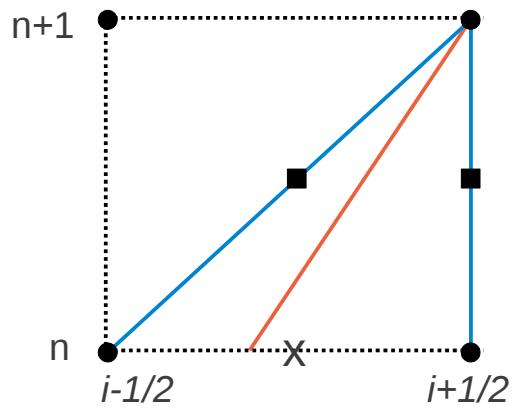


Interpolation of flux variables

$$\rho_{i+\frac{1}{2}}^{n+\frac{1}{2}} + \rho_{i-\frac{1}{2}}^{n-\frac{1}{2}} = 2\rho_i^n$$

CABARET monotonization procedure

Goloviznin et al., 2003,,



$$\frac{d \rho}{dt} = f(x, t)$$

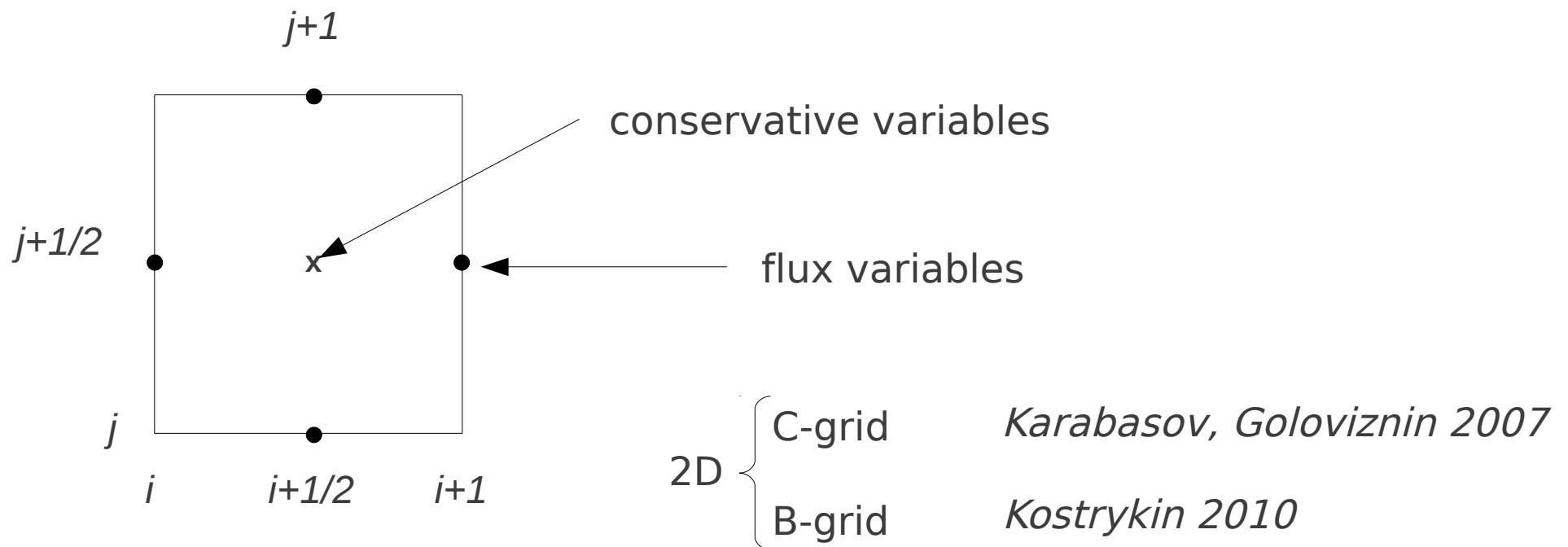
$$\rho_{i+1/2}^{n+1} = \rho_o^n + \int_{t^n}^{t^{n+1}} f(x(t), t) dt \approx \rho_o^n + f(x(t^{n+1/2}), t^{n+1/2}) \tau$$

$$m_i^n = \rho_{i-1/2}^n + f(x_i, t^{n+1/2}) \tau, M_i^n = \rho_{i+1/2}^n + f(x_{i+1/2}, t^{n+1/2}) \tau$$

$$\min(m_i, M_i) \leq \rho_{i+1/2}^{n+1} \leq \max(m_i, M_i)$$

CABARET advection scheme in 2D case

$$\frac{\partial \rho}{\partial t} + \frac{\partial u \rho}{\partial x} + \frac{\partial v \rho}{\partial y} = 0$$



CABARET advection scheme in 2D case

I. Update conservative variables

$$\rho_{i+1/2, j+1/2}^{n+1/2} = \rho_{i+1/2, j+1/2}^{n-1/2} - \tilde{u}(\rho_{i+1, j+1/2}^n - \rho_{i, j+1/2}^n) - \tilde{v}(\rho_{i+1/2, j+1}^n - \rho_{i+1/2, j}^n)$$

II. Update flux variables

$$\hat{\rho}_{i+1, j+1/2}^{n+1} = 2\rho_{i+1/2, j+1/2}^{n+1/2} - \rho_{i, j+1/2}^n, \hat{\rho}_{i+1/2, j+1}^{n+1} = 2\rho_{i+1/2, j+1/2}^{n+1/2} - \rho_{i+1/2, j}^n$$

III. Flux variables correction

$$\begin{aligned} m_{i+1, j+1/2} &= \min(\rho_{i+1, j+1/2}^n, \rho_{i, j+1/2}^n), M_{i+1, j+1/2} = \max(\rho_{i+1, j+1/2}^n, \rho_{i, j+1/2}^n), \\ m_{i+1/2, j+1} &= \min(\rho_{i+1/2, j+1}^n, \rho_{i+1/2, j}^n), M_{i+1/2, j+1} = \max(\rho_{i+1/2, j+1}^n, \rho_{i+1/2, j}^n), \\ \rho_{i+1, j+1/2}^{n+1} &= \max(\min(\hat{\rho}_{i+1, j+1/2}^{n+1}, M_{i+1, j+1/2}), m_{i+1, j+1/2}), \\ \rho_{i+1/2, j+1}^{n+1} &= \max(\min(\hat{\rho}_{i+1/2, j+1}^{n+1}, M_{i+1/2, j+1}), m_{i+1/2, j+1}) \end{aligned}$$

Monotonicity of the CABARET scheme

1D *Ostapenko, 2009*

2D,3D ?

IIIa

$$\begin{aligned} m_{i+1,j+1/2} &= \min(\rho_{i+1,j+1/2}^n, \rho_{i,j+1/2}^n), M_{i+1,j+1/2} = \max(\rho_{i+1,j+1/2}^n, \rho_{i,j+1/2}^n), \\ m_{i+1/2,j+1} &= \min(\rho_{i+1/2,j+1}^n, \rho_{i+1/2,j}^n), M_{i+1/2,j+1} = \max(\rho_{i+1/2,j+1}^n, \rho_{i+1/2,j}^n), \\ \rho_{i+1,j+1/2}^{n+1} &= \max(0., \max(\min(\hat{\rho}_{i+1,j+1/2}^{n+1}, M_{i+1,j+1/2}), m_{i+1,j+1/2})), \\ \rho_{i+1/2,j+1}^{n+1} &= \max(0., \max(\min(\hat{\rho}_{i+1/2,j+1}^{n+1}, M_{i+1/2,j+1}), m_{i+1/2,j+1})) \end{aligned}$$

Statement. Under condition $\tilde{u} \leq \frac{1}{4}, \tilde{v} \leq \frac{1}{4}$ scheme I, II, IIIa is positive.

Proof

$$\rho_{i+1/2,j+1/2}^{n+1/2} \geq \rho_{i+1/2,j+1/2}^{n-1/2} - u \rho_{i+1,j+1/2}^n - v \rho_{i+1/2,j+1}^n$$

$$\rho_{i+1/2,j+1/2}^{n-1/2} = \frac{1}{2} (\rho_{i,j+1/2}^{n-1} + \rho_{i+1,j+1/2}^n),$$

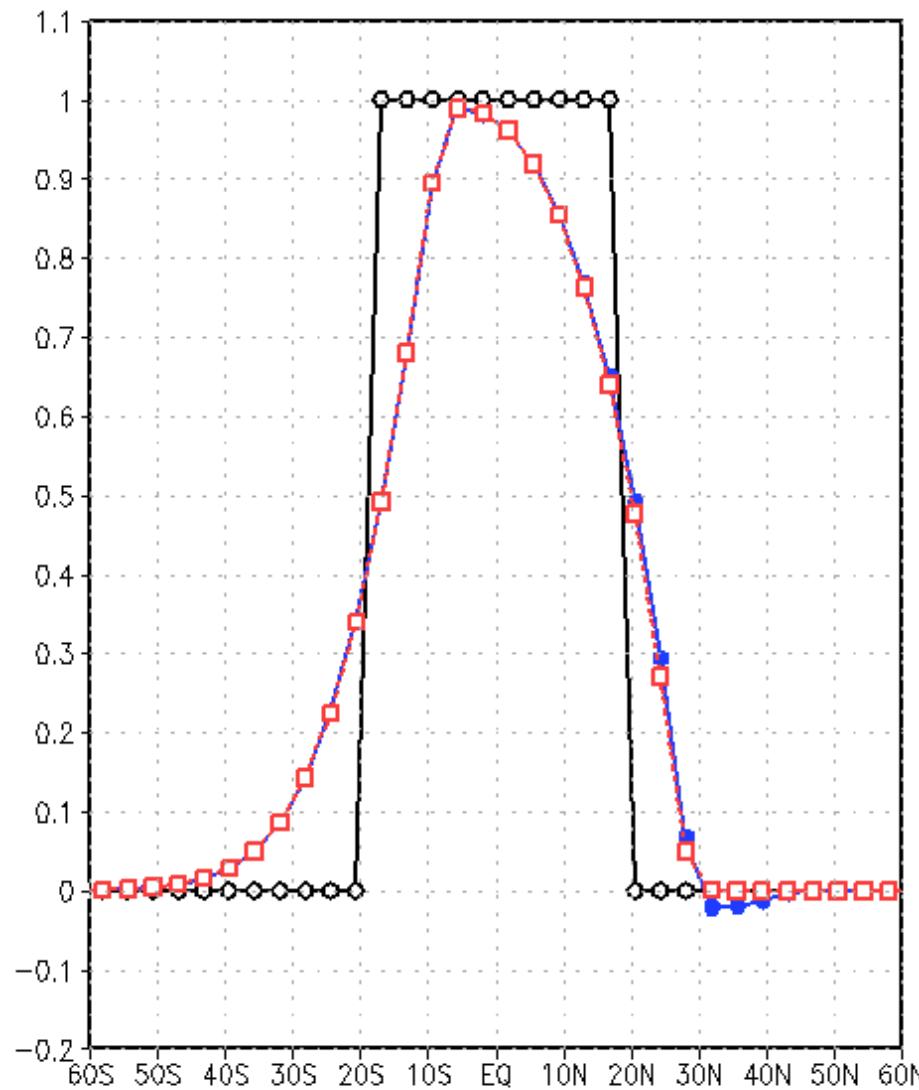
$$\rho_{i+1/2,j+1/2}^{n-1/2} = \frac{1}{2} (\rho_{i+1/2,j}^{n-1} + \rho_{i+1/2,j+1}^n),$$

$$\rho_{i+1/2,j+1/2}^{n+1/2} \geq \left(\frac{1}{4} - \tilde{u}\right) \rho_{i+1,j+1/2}^n + \tilde{u} \rho_{i,j+1/2}^{n-1} + \left(\frac{1}{4} - \tilde{v}\right) \rho_{i+1/2,j+1}^n + \tilde{v} \rho_{i+1/2,j}^{n-1}.$$

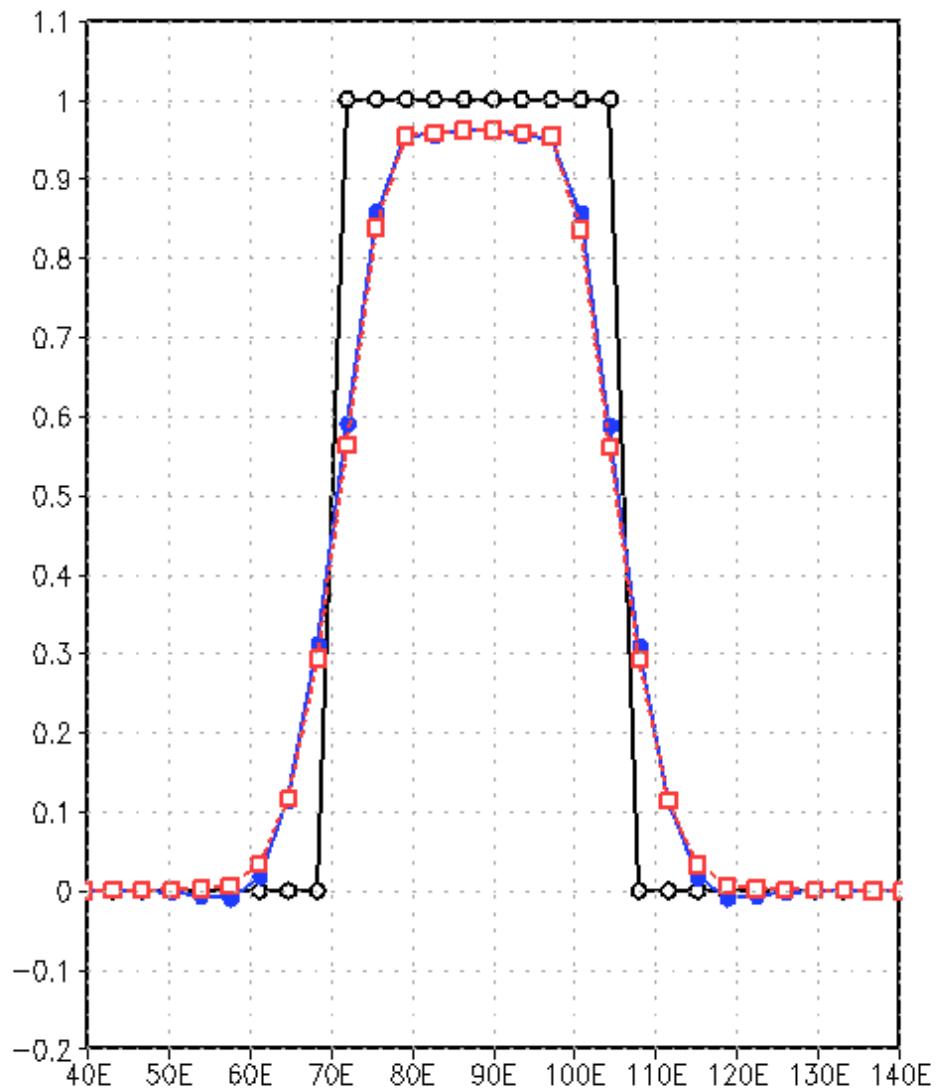
Results of the solid-body rotation test on the sphere

$$\alpha = \frac{\pi}{2}$$

S-N



E-W



Problem of pressure-tracer inconsistency

$$\frac{\partial \pi q}{\partial t} + \operatorname{div}(\pi q \mathbf{u}) = 0, \quad (1)$$

$$\frac{\partial \pi}{\partial t} + \operatorname{div}(\pi \mathbf{u}) = 0, \quad (2)$$

$$\frac{\partial \pi}{\partial t} + \int_0^1 \operatorname{div}_2(\pi \mathbf{u}) d\sigma = 0 \quad (3)$$

$$\frac{(\pi q)^{n+1} - (\pi q)^n}{\Delta t} + F^{n+1/2}(\pi q \mathbf{u}) = 0, \quad (4)$$

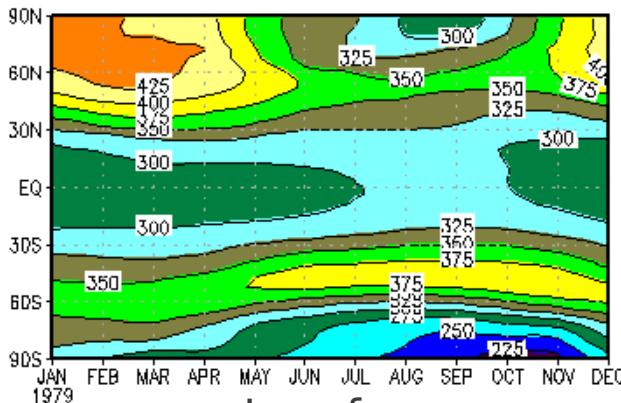
$$\frac{\pi^{n+1} - \pi^n}{\Delta t} + F_2^{n+1/2}(\pi \mathbf{u}) + (\pi \dot{\sigma})_\sigma^{n+1/2} = 0, \quad (5)$$

$$\frac{\pi^{n+1} - \pi^n}{\Delta t} + \sum_{i=1}^N F_2^{n+1/2}(\pi \mathbf{u}_i) = 0. \quad (6)$$

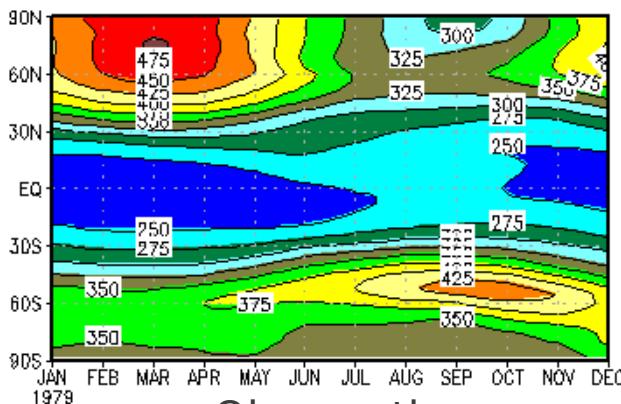
- 1) For global conservation one should know the pressure field on the next time step before advection
- 2) Inconsistency in calculation of advection of tracer and pressure could lead to the **non-monotonicity** of the tracer distribution

Seasonal cycle of total ozone column

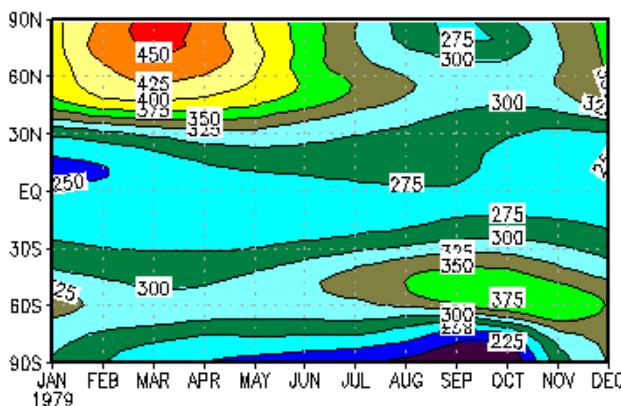
CABARET



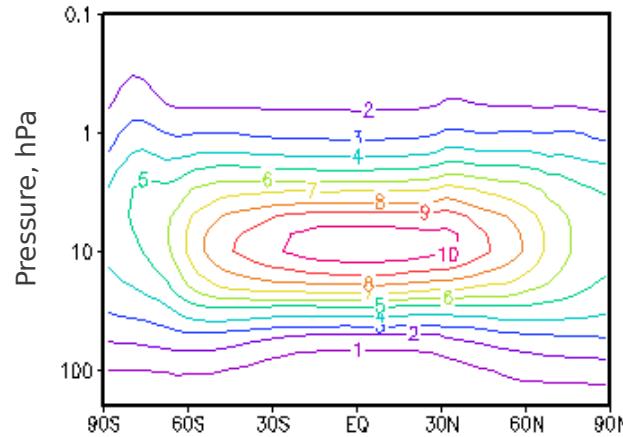
Leapfrog



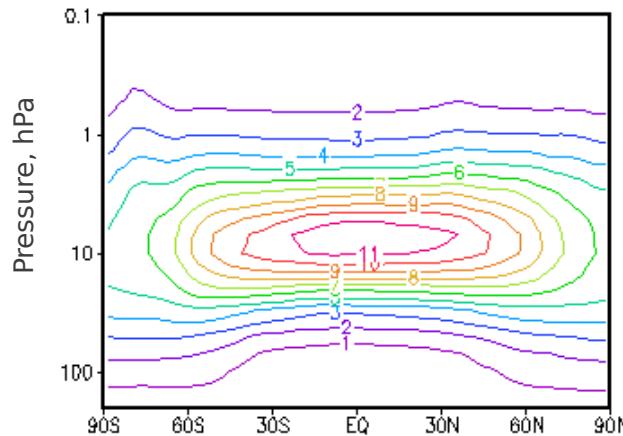
Observations



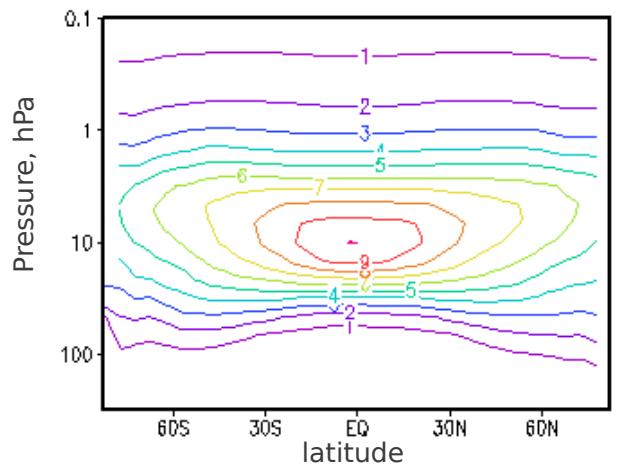
Annual distribution of Ozone (ppmv)



CABARET

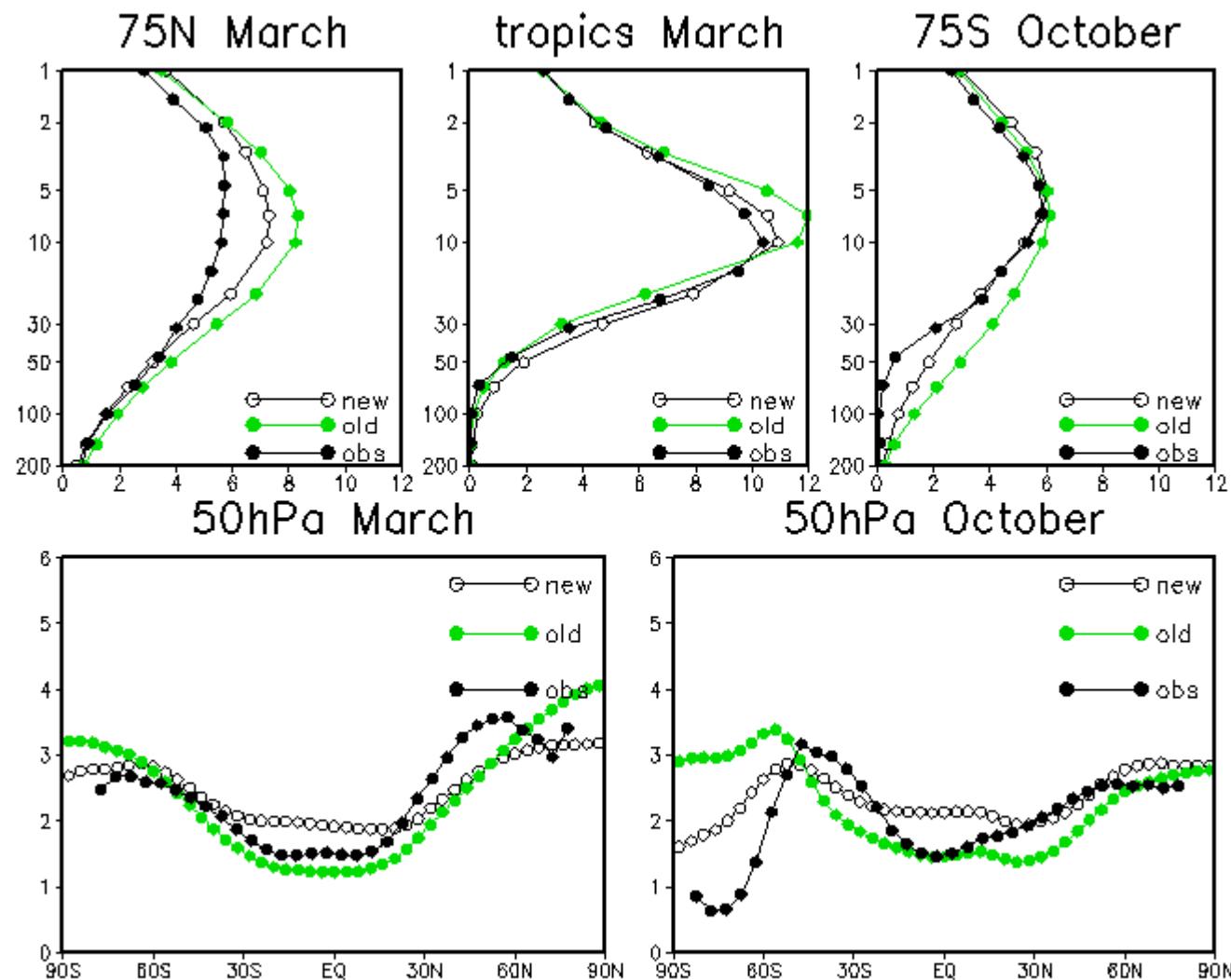


leapfrog

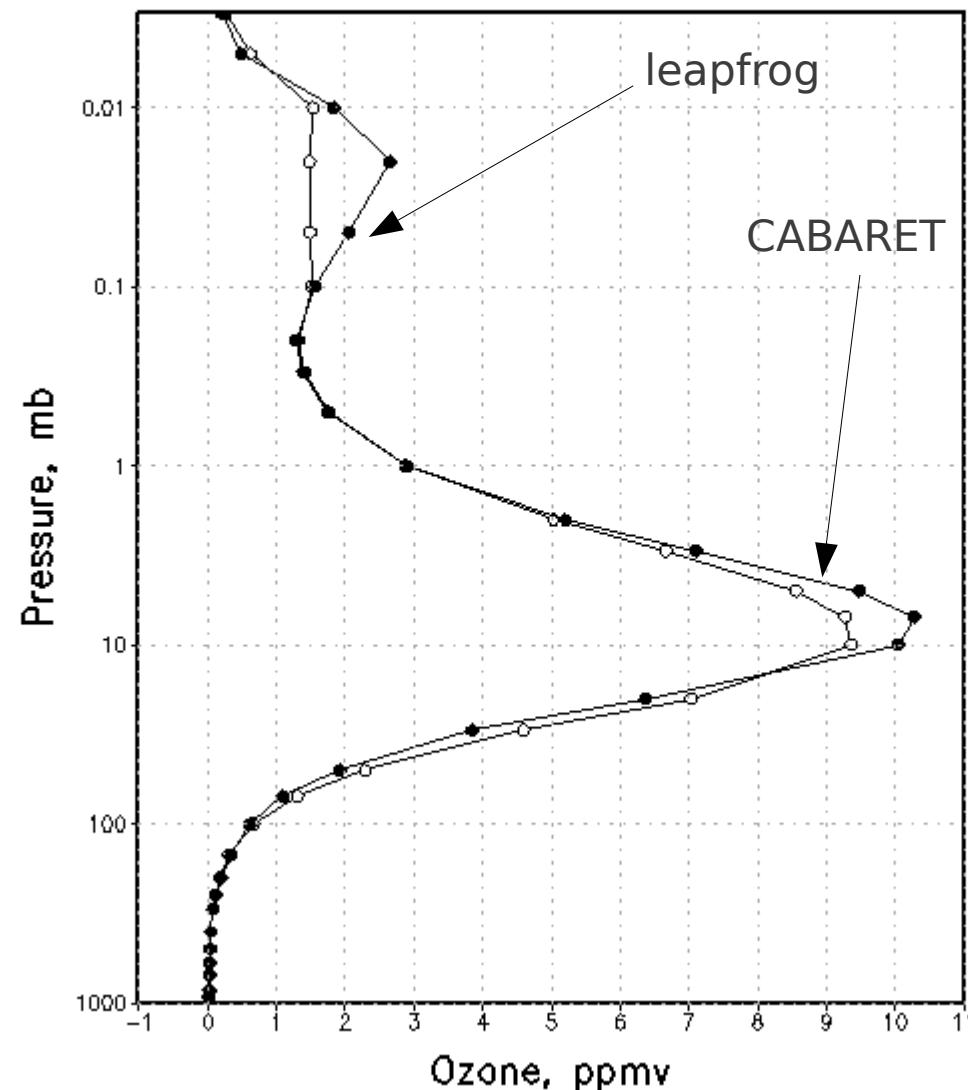


HALOE

Vertical profiles ozone



Zonally mean annual ozone profile

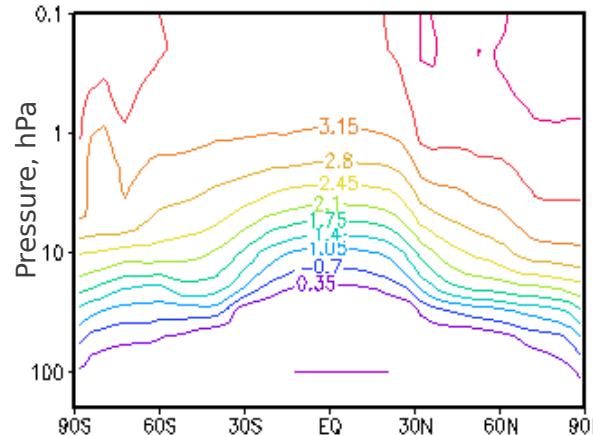


Conclusions:

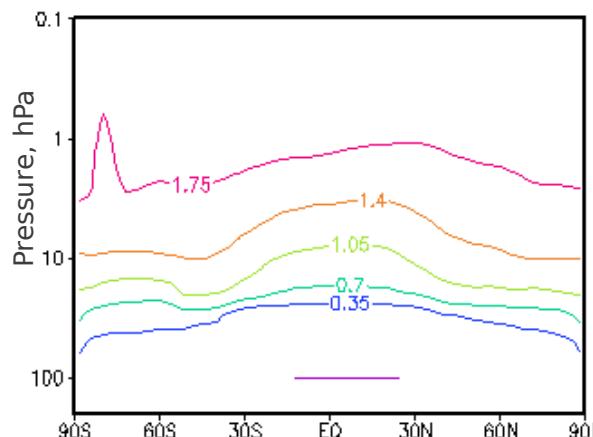
- 1) A positive multidimensional version of CABARET scheme is proposed
- 2) The implementation of the new advection schemes into chemical transport model leads to:
 - improvement of ozone climatology near stratopause (disappear artificial maximum) and near stratospheric ozone maxima (decreasing)
 - improvement in the description of total ozone column at high and middle latitudes SH (deeper ozone holes, weaker midlatitude maxima)
 - at tropics a total ozone column becomes larger than in observations
- 3) Elapsed time for CTM increases on 30%. Memory storage increases in 4 times for each tracer.

Thanks for your attention!

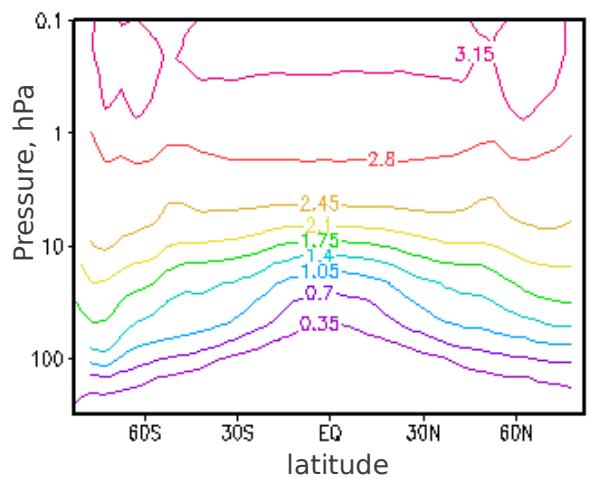
Annual distribution of HCl (ppbv)



CABARET



leapfrog



HALOE