

Eddy diffusion in the atmosphere and at the ocean surface

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$k = 0.17 \text{cm}^2/\text{s}$ diffusion coefficient of O_2 into N_2
diffusion equation, Adolf Fick, 1855,
after Charles Fourier, 1821.

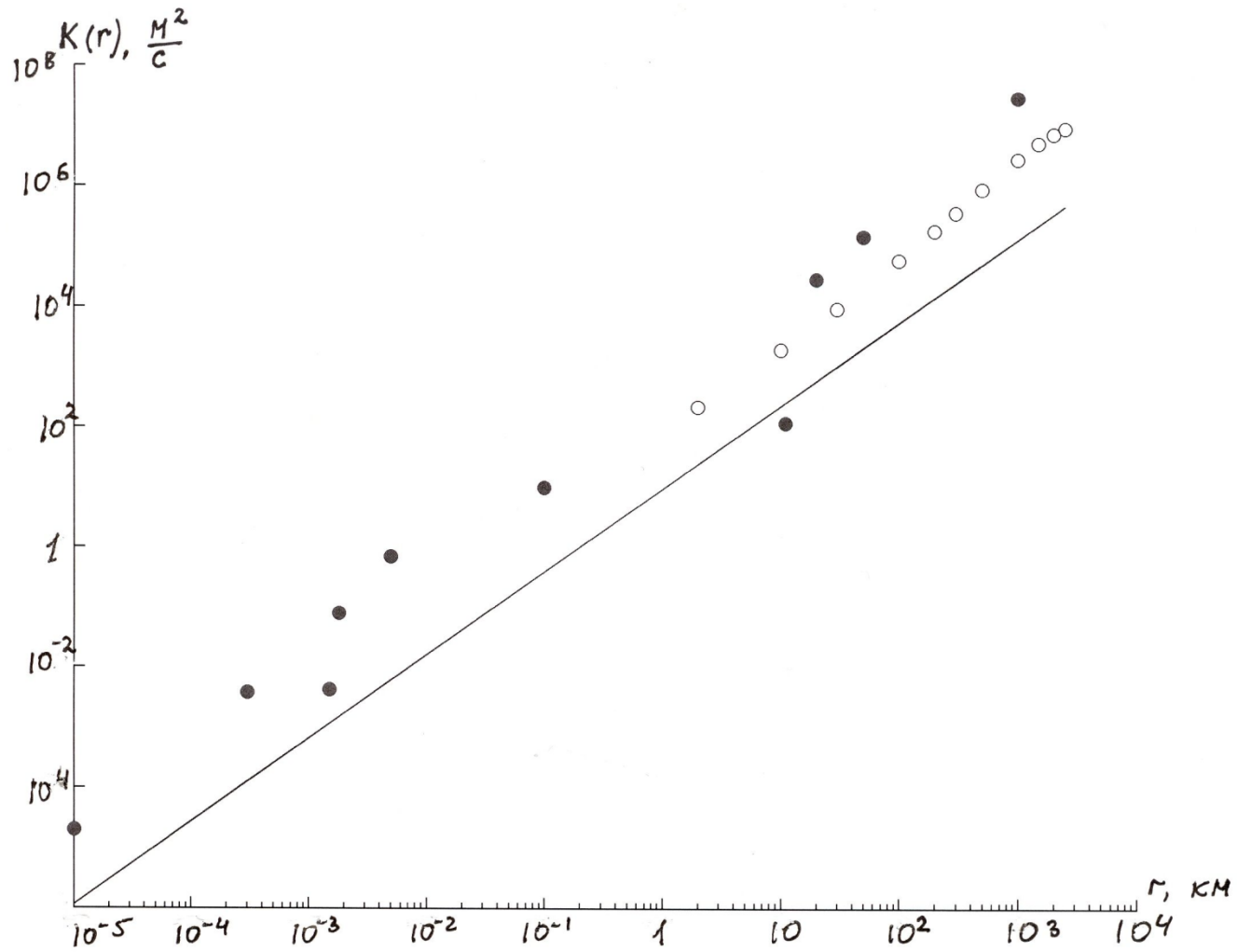
$K \sim r u$, eddy diffusion coefficient,
G.I. Taylor. 1915.

$K \sim r^{4/3}$ for the atmosphere,

L.F. Richardson 1926, 1929.

Compare the last two lines and get $u \sim r^{1/3}$!

Kolmogorov 1941



- Richardson 1926, 1929
- Golitsyn 2001

L.F. Richardson and Stommel, 6 January, 1948

(J. Meteorology 1948. V. 5. No. 5. 238 – 240)

parsnip white pieces ~ 1 inch relative distance l with time t

$$K(l) = \frac{1}{n} \frac{(l_0 - n^{-1} l)^2}{2t}$$

l_0 - initial distance between markers at $t = t_0$

~100 markers for about an hour, for $l = 3\text{m}$, $t = 30$ sec.

$$K(l) \sim l^n, \quad n = 1.4 \quad \text{in ref. 1948}$$

recalculation by myself $n = 1.32$.

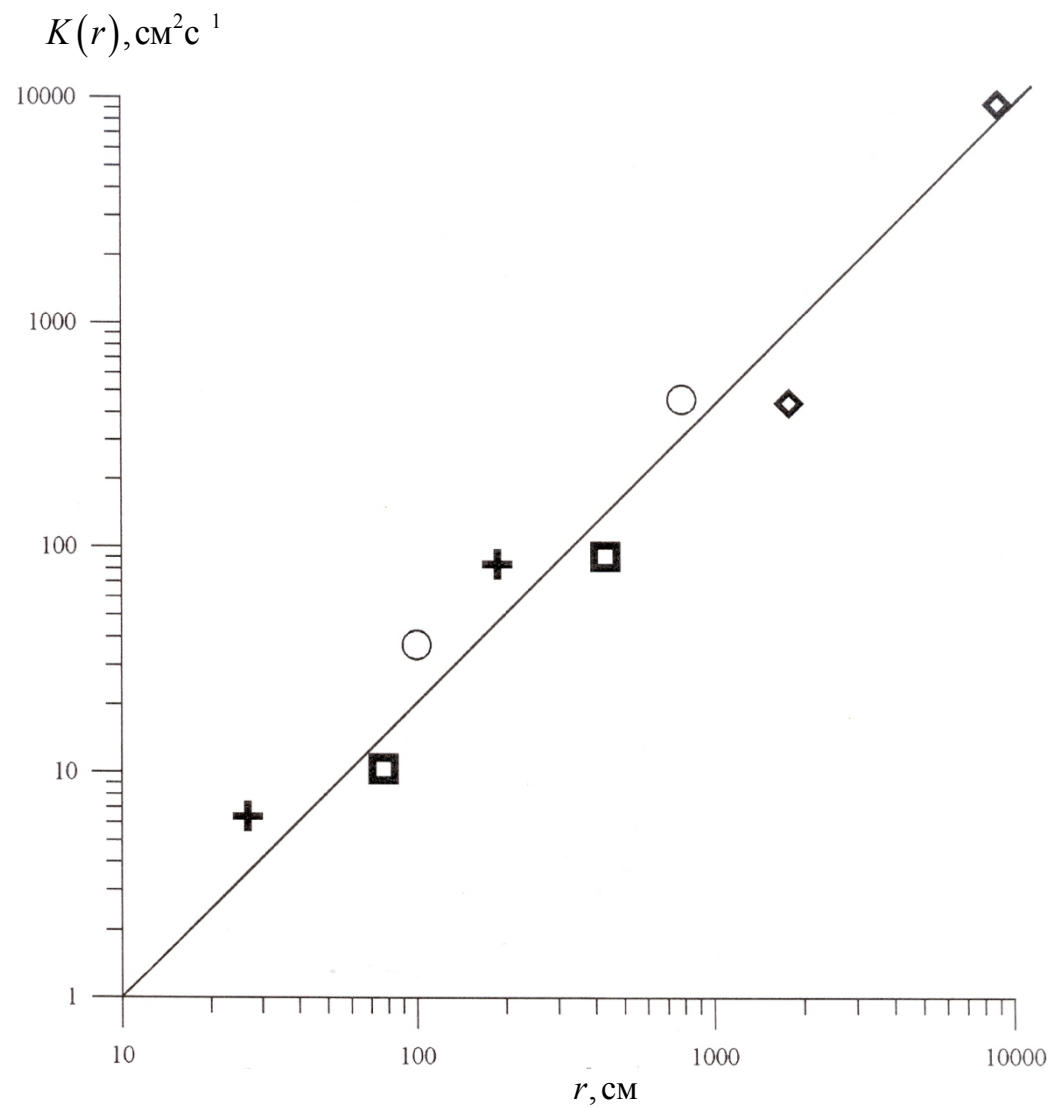
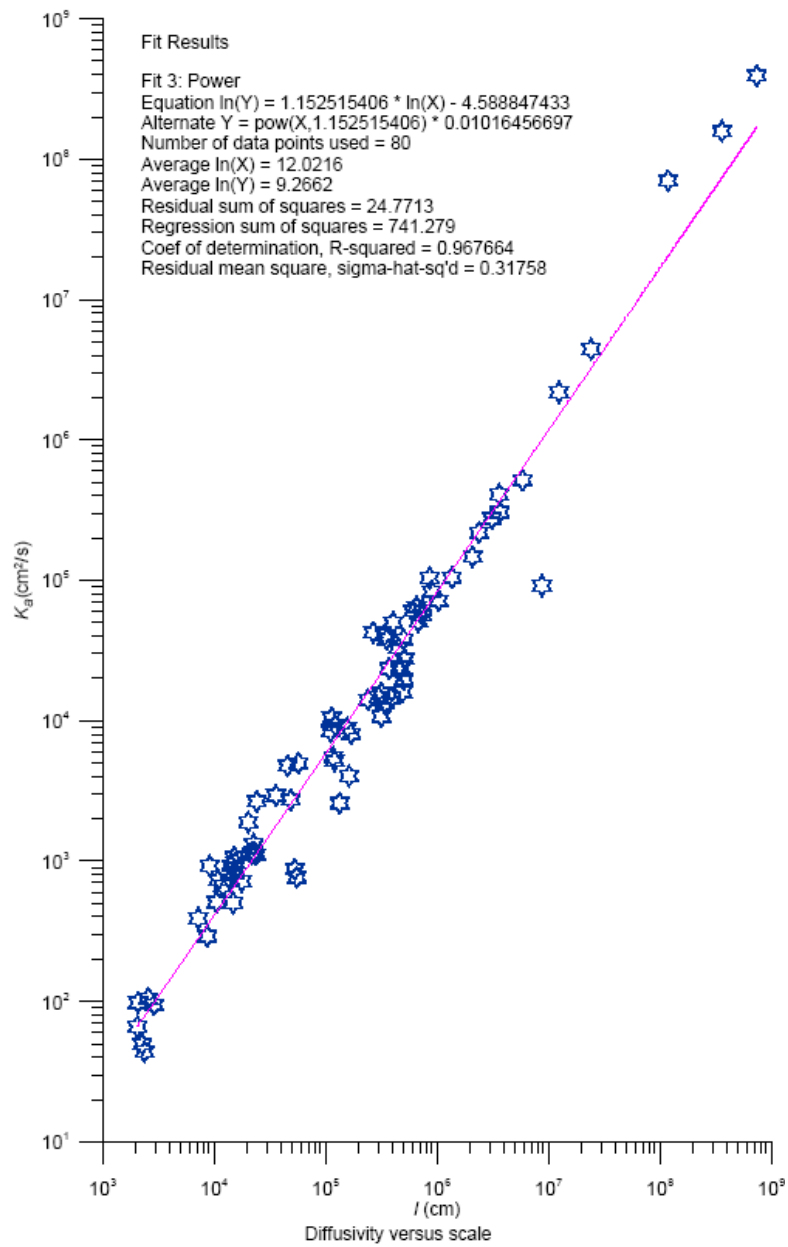
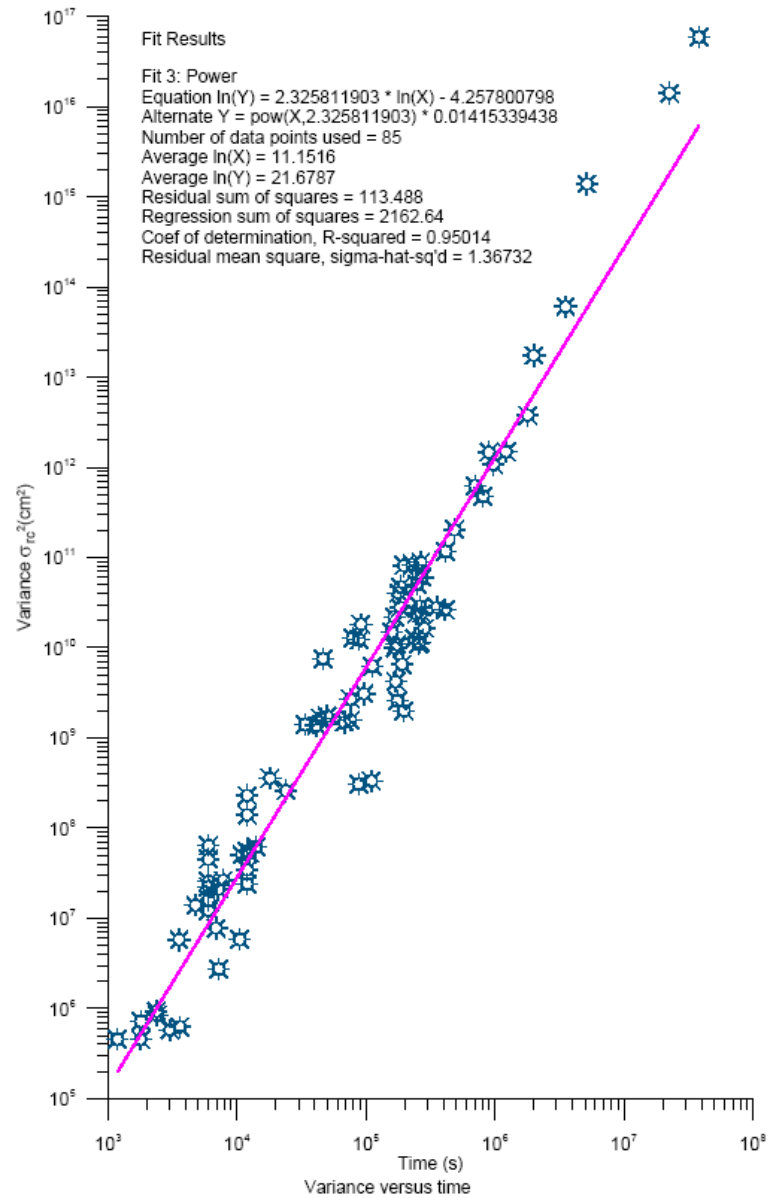
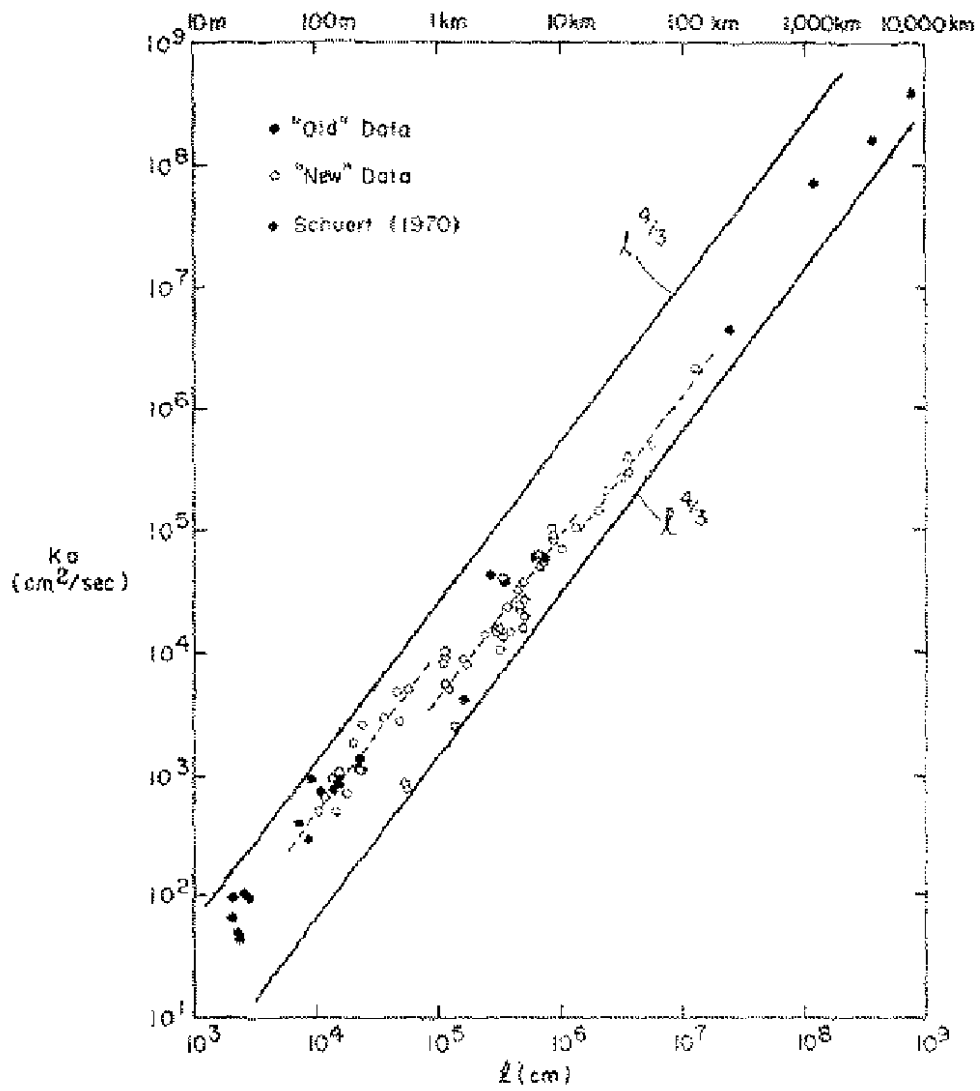


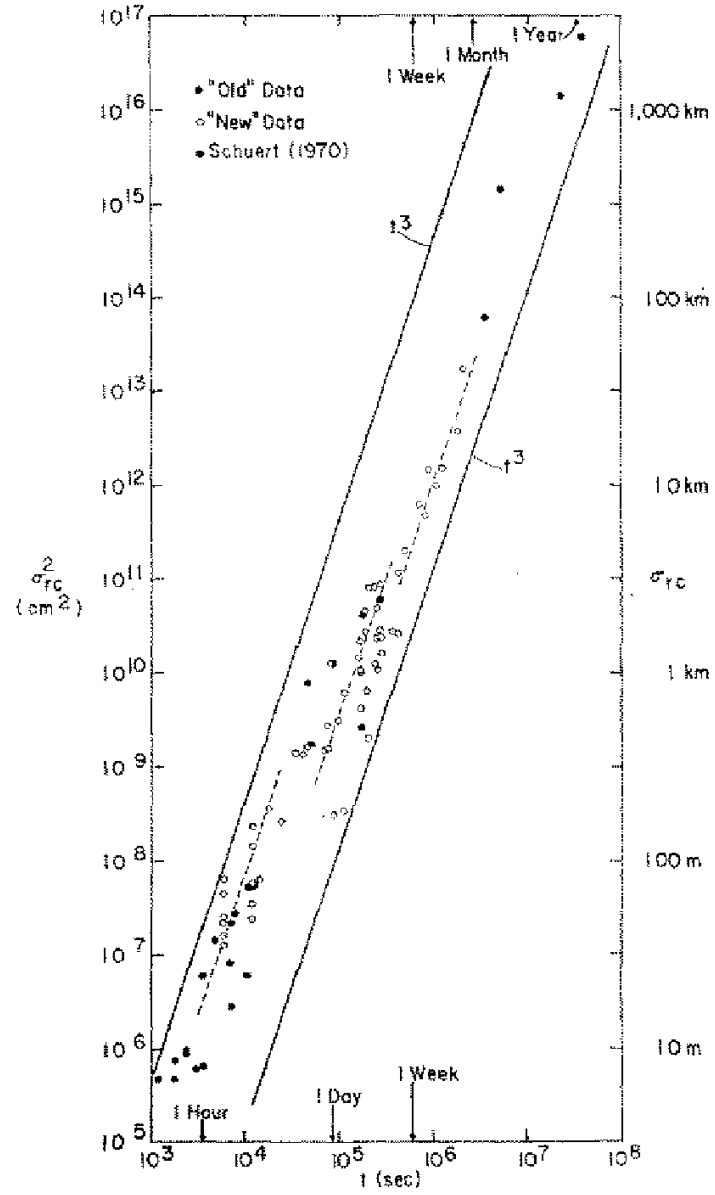
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3 c



Taylor $K(r) \sim r^{-5/3}$, $u = \left\langle D_u(r)^{1/2} \right\rangle = \left\{ \left\langle u(x+r) - u(x) \right\rangle^2 \right\}^{1/2}$

how to get $D_u(r)$, the structure function for velocity field in the sea surface waves?

We know the time elevation spectrum of the surface

$$E_z(\omega) \sim \omega^{-n}, \quad n(\omega) = \frac{U_{10}}{c} \quad \text{- the wave age}$$

vertical velocity: $w = \frac{dz}{dt}, \quad E_w(\omega) = \omega^{-2} E_z(\omega).$

Due to water incompressibility $E_u(\omega) \sim E_w(\omega) = \omega^2 E_z(\omega)$

dispersion relation: $\omega^2 = kg$

(Golitsyn 2007 general case for $n = 4$)

Probability transformation $E(\omega) d\omega = E(k) dk$

$$E_u(k) = c_{gr} E_u(\omega) \sim c_{gr} kg E_z(k)$$

$$D_u(r) = 2 \int_0^r (1 - \cos r) E_u(k) dk$$

$$D_u(r) \sim r^{\frac{n+3}{2}}; \quad K(r) \sim r^{\frac{n+1}{4}}$$

From the parabolic nature of the diffusion equation with scale variable
diffusion coefficient $K(r) : r$

$$r^2 \quad S(t) \sim t^{\frac{2}{n}} = t, \quad = \frac{8}{7} \quad \text{at} \quad = \frac{n+1}{4}$$

Okubo $= 1.15$ for $r = 1000\text{km}$

$= 2.34$ for $1 \text{ day} < t < 1 \text{ month}$

3 locally

$n = 4$ Kitaigorodsky 1962, Zakharov 1966, Toba 1973

Energy transfer: $D_u(r) \sim r^{1/2}$; $K(r) \sim r^{5/4}$, $\alpha = 1.25$, $S(t) \sim t^{8/3}$

$n = \frac{13}{3} = 4 + \frac{1}{3}$; Hasselmann K., Hasselmann S. 1974

Spectral momentum transfer

$$D_u(r) \sim r^{2/3}, \quad K(r) \sim r^{4/3}, \quad S(t) \sim t^3,$$

$n = \frac{11}{3} = 4 - \frac{1}{3}$; Zakharov&Zaslavsky 1983:

Spectral action transfer: $D_u(r) \sim r^{1/3}$, $K(r) \sim r^{7/6}$, $S(t) \sim t^{12/5}$,

$$\sim t^{-1}, \quad \gamma_1 = 1/2, \quad \text{Golitsyn 2010}$$

Gagnaire-Renou E., Benoit M., Badulin S.I. J. Fluid Mech. 2011
 computations:

$$n = \frac{13}{3} \quad \text{for} \quad \gamma_1 > 2, \quad \text{young waves, small fetch, } t \sim 1.5\text{hr}$$

$$= 4/3, \quad \gamma_1 = 3 \quad \text{Richardson \& Stommel: } \gamma_1 = 4/3!$$

$$n = 4 \quad \text{for} \quad 1.2 < \gamma_1 < 2 \quad = 5/4, \quad \gamma_1 = 8/3, \quad t = 2 \quad 4\text{hr}$$

no reliable data

$$n = 11/3 \quad \text{for} \quad 0.83 < \gamma_1 < 1.2 \quad \text{old waves near saturation } t = 4\text{hr}$$

$$= 7/6, \quad \gamma_1 = 2.4 \quad \text{Okubo: } \gamma_1 = 1.15, \quad \gamma_1 = 2.34!$$

Horace Lamb, 1895. Hydrodynamics

§ 349. Surface waves with viscosity. Linear approximation

$$u = (ikAe^{kz} + mCe^{mz})e^{ikx+nt}, \quad i = e^{i/2}$$

$$w = (kAe^{kz} - ikCe^{mz})e^{ikx+nt}, \quad z < 0$$

Small parameter

$$\epsilon = k^2 l^2, \quad k = 2\pi / \lambda,$$

$$n = 2k^2 \pm i = (2\epsilon \pm i)$$

$$\frac{C}{A} = m \frac{2k^2}{n} = m \frac{2\epsilon}{2\epsilon \pm i} \ll 1$$

$$m^2 = k^2 + n$$

$$m^2 = k^2 \pm i$$

The depth of the vorticity viscous layer $l \sim (\nu / \omega)^{1/2}$

$m^2 = k^2 \pm i / \dots$ complex number with large imaginary part on the

complex plane:

$$m^2 = \left(k^4 + \dots / \dots \right)^{1/2} e^i, \quad = \arctg \frac{m}{k^2} = \arctg(\pm 1 / \dots) = \frac{\dots}{2} \dots,$$

$$m = \left(k^4 + \dots / \dots \right)^{1/4} e^{i/2} = \frac{k}{1^{1/2}} e^{i(\dots/4)}.$$

For typical ocean waves (Golitsyn, 2010) $h = 2.7$ m, $\dots = 40$ m, $\dots = 1.25 \text{ c}^{-1}$,

$$\dots_1 = \frac{k^2}{\dots} = \frac{4 \dots^2}{2} = 2.4 \cdot 10^8, \quad \dots_1^{1/2} = 1.56 \cdot 10^4$$

At $z = 0$:

$$u = (ikA + mC)e = kA \left(i + \frac{m c}{k A} \right) e = kA \left(i + (2)^{1/2} \right) e$$

$$w = (kA - ikC)e, \quad = ikx + nt = i(kx \pm t) - 2k^2t = i - 2k^2t$$

$$u = \left(1 + (2)^{1/2} \right)^{1/2} e^{i\phi}, \quad \phi = \arctg \frac{1}{(2)^{1/2}} = 90^\circ, \quad \phi = 0.69 = 41^\circ$$

$$w = kA(1 - ic/A)e = kA(1 - 2i)e^{i\phi} = kA(1 + 4)^{1/3} e^{i\phi},$$

$$\phi = \arctg 2 = 90^\circ, \quad \phi = 0.004$$

permanent addition to the phase of horizontal wave of order one minute!

Comparing to the phase of vertical component!

Slow horizontal diffusion motion on x !

Mean momentum in time:

$$\langle uw \rangle = \left\langle k^2 A^2 \sin\left(\omega_1 + 2\omega_1^{1/2}\right) \cos\left(\omega_1 + 2\omega_1^{1/2}\right) \right\rangle = k^2 A^2 \left\langle \left(\sin\left(\omega_1 + 2\omega_1^{1/2}\right) \cos\left(\omega_1 + 2\omega_1^{1/2}\right) \pm \cos\left(\omega_1 + 2\omega_1^{1/2}\right) \sin\left(\omega_1 + 2\omega_1^{1/2}\right) \right) \cos\left(\omega_1 + 2\omega_1^{1/2}\right) \right\rangle$$

$$k^2 A^2 \omega_1^{1/2} = \frac{1}{2} h^2 \omega_1^{1/2} = \frac{1}{2} h^2 \omega_1^{1/2}.$$

For the mean waves this corresponds to 1 cm/s

For the mean wind of 8 m/s the Stokes drift:

$$u_a = 29 \text{ cm/s and } u_{\text{drift}} = 0.5u_a = 15 \text{ cm/s, J. Wu 1975.}$$

But the non-linear Stokes drift is over the whole region of wind action moving as a whole the pollution spot, synoptic scale 1000 km.

$$r^2 = 0.0108t^{2.34}, \quad r = 0.1t^{1.17} \quad \text{slightly overballistic motion of a spot boundary}$$

$$r^2 = t^3 \quad \text{for isotropic turbulence } r \sim t^{1.5}, \quad \text{when } K \sim \frac{1}{3} r^{4/3}$$

Due to G.I. Taylor (1915) the diffusion coefficient

$$K(r) = \langle a_2 r^{-1/4} u(r) \rangle = a_2 r^{-1/4} D_u(r)^{1/2}, \quad a_2 \text{ number.}$$

For $n = 11/3$:

$$K(r) = a_2 r^{-1/4} \left(3.52 \frac{8}{3} h^2 \frac{8/3}{p} g^{1/3} r^{1/3} \right)^{1/2} \sim r^{7/6}$$

Okubo: $K(r) = 0.0103 r^{1.15}$

We recalculated from the tables by Okubo $n = 1.15 \pm 0.05$

Comparing the theory with experiment for the annual mean wave field for the

World Ocean (Golitsyn, 2010) we find $a_2^{-1/2} = 1.56 \cdot 10^4$ and $a_2 = 2.3 \cdot 10^{-3}$.

Okubo for the area of tracer spot:

$$S(t) = 0.0108t^{2.34} = r^2, \quad \text{in cm}^2.$$

Diffusion equation

$$\frac{S}{t} = \frac{K(r)}{r} \frac{S}{r}, \quad K(r) = br, \quad ,$$

where $[b] = L^2 T^{-1}$ - constant over t and r . Introduce $\tau = bt$, $[r] = L$.
 Dimensional analysis gives for $[S] = r^2$:

$$S = b_1 r^2, \quad \tau = \frac{2}{2.34}$$

With $\tau = 7/6$ we get $\tau = 12/5 = 2.4$, while reanalyzing the Okubo tables

$\tau = 2.33 \pm 0.10$. Our result gives for $b = 460$. If we take 1971 value $\tau = 1.15$

then $\tau = 2(2.34)^{-1} = 2.35$. Both values $\tau = 1.15$ and $\tau = 2.34$

Okubo obtained by eye!!

From Stommel $K(r) = 4.6 \cdot 10^{-2} r^{4/3}$, our estimate

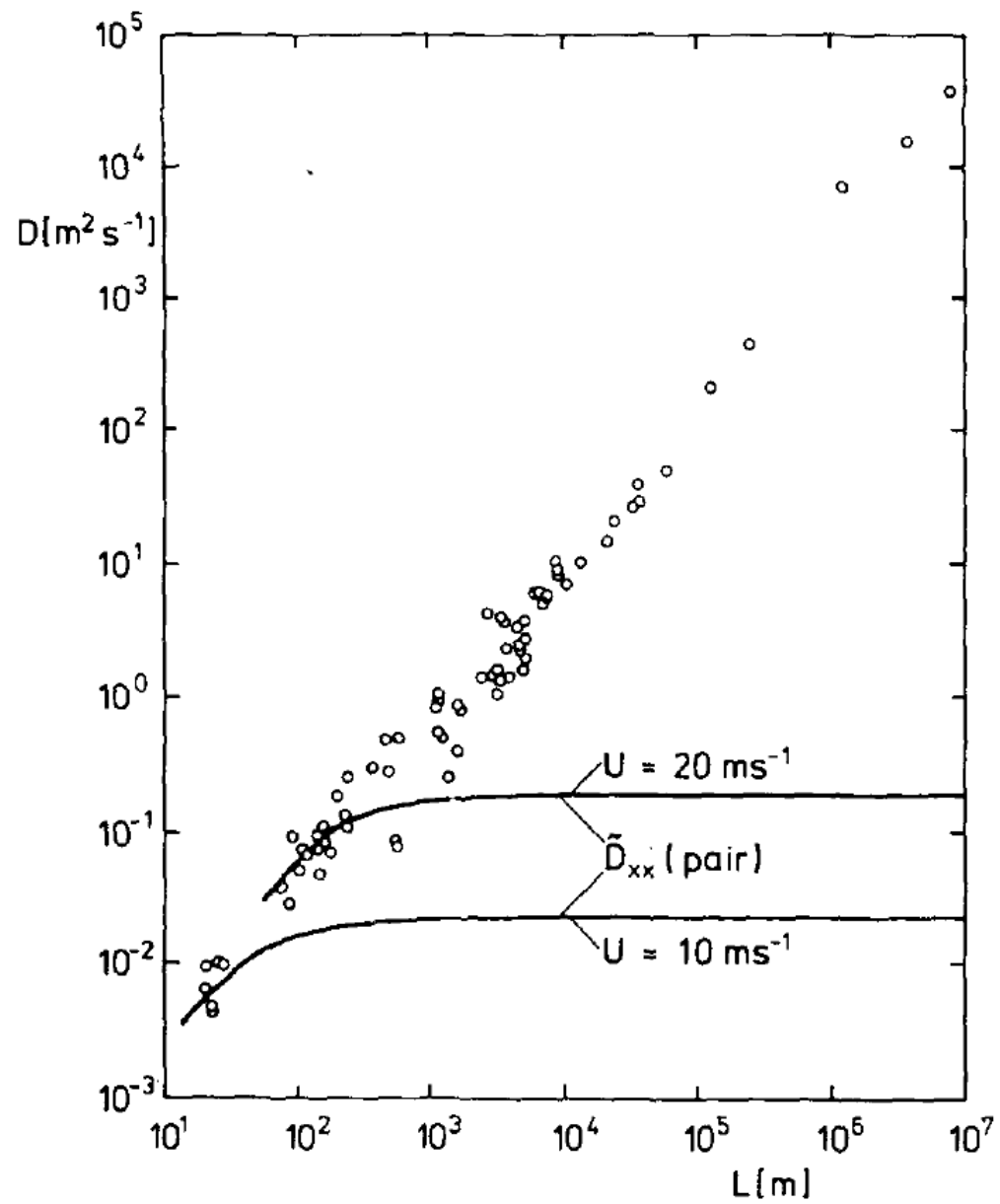
for $\frac{2}{2-} = 3$, i.e. $S = r^2 \sim t^3$,

as for the case of locally homogeneous and isotropic turbulence, Batchelor, 1950.

This theory based on linear approximation (Lamb, 1895) is self-consistent but determines constants, which depend on wind and fetch, from observations.

The limits of observational results corresponds to theory and numerics of Gagnaire-Renou E.

Benoit M., Badulin S.I. 2011, JFM, 669, 178-213.



K. Herterich & K. Hasselmann

The horizontal diffusion of tracers by surface waves JPO
1982. V. 12, 704 – 711.

Random fluctuations of the local Stokes-drift current,
Pirson_Moscowitz can be explained wave spectrum diffusion
coefficients for single particle, particle pairs and continuous
traces spot small scales up to hundreds m .

Thanks for your attention!